

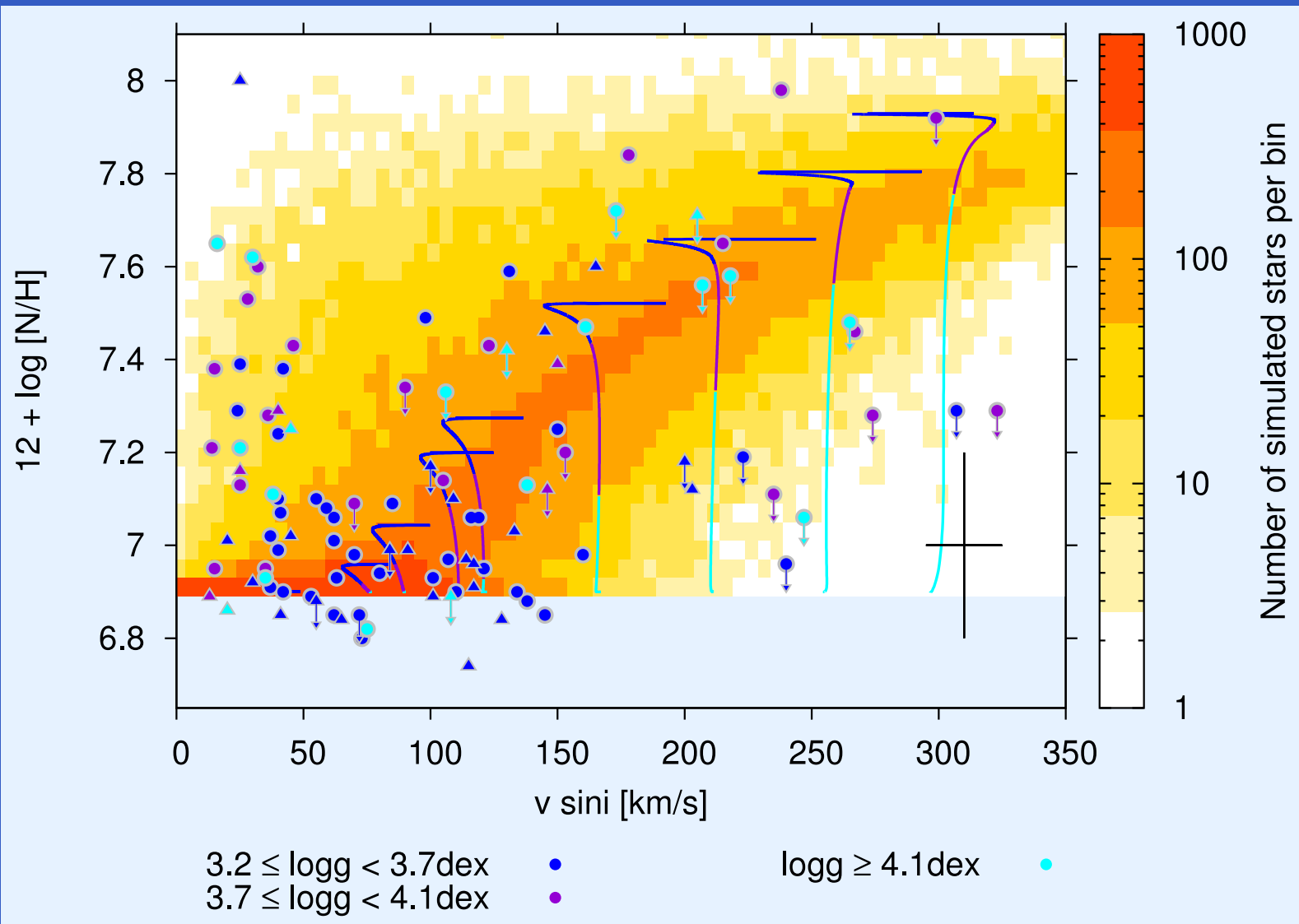
Very Massive Stars

Norbert Langer (Bonn University)

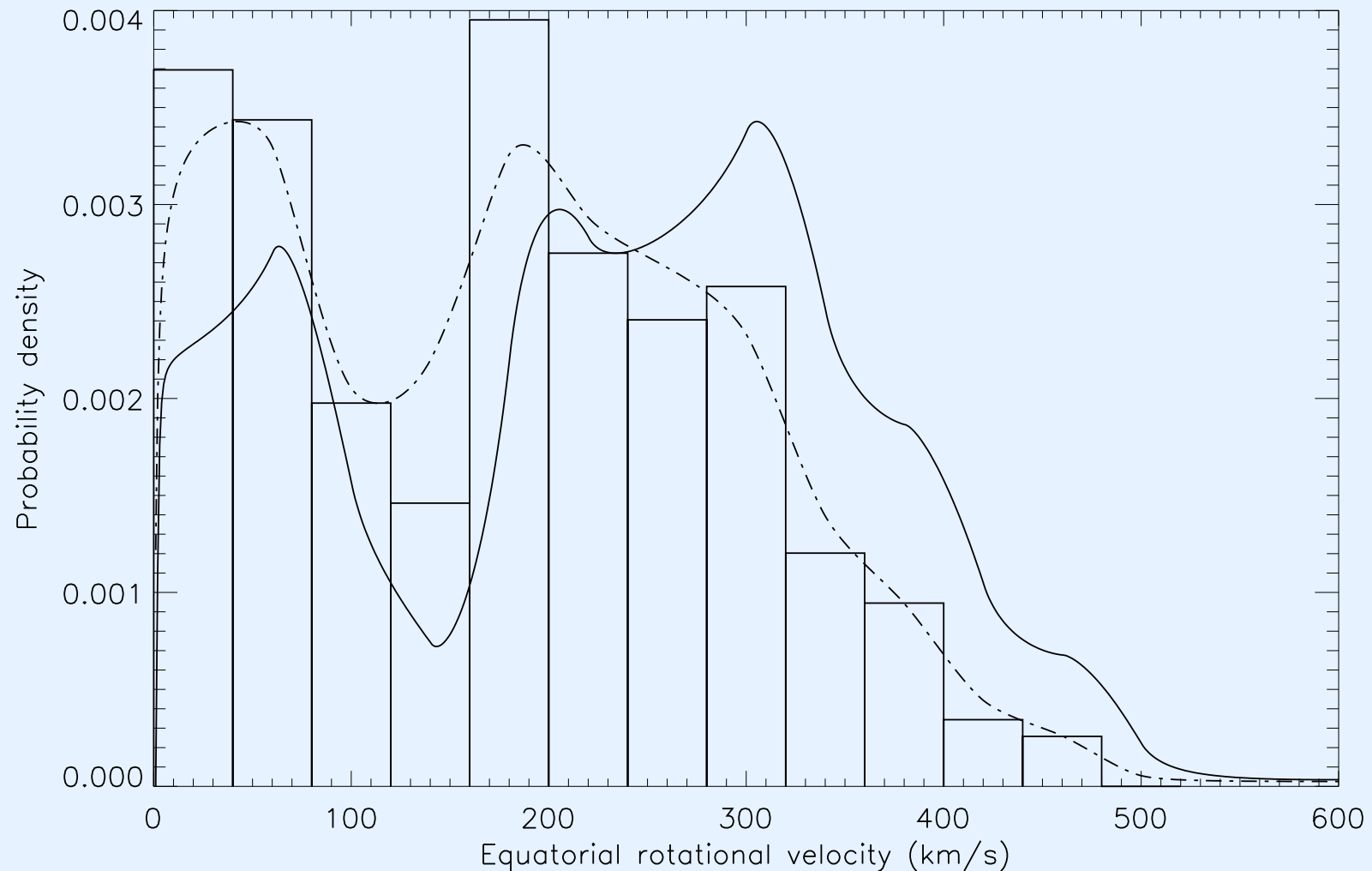
with

- Ines Brott (Vienna)
- Matteo Cantiello (UCSB)
- Selma de Mink (STScI)
- Karen Köhler (Bonn)
- Debashis Sanyal (Bonn)
- Alexandra Kozyreva (Bonn)
- Colin Norman (JHU)
- Alex Heger (Monash)
- Alex de Koter (Amsterdam)
- Onno Pols (Nijmegen)
- Jorick Vink (Armagh)
- Sung-Chul Yoon (Bonn)

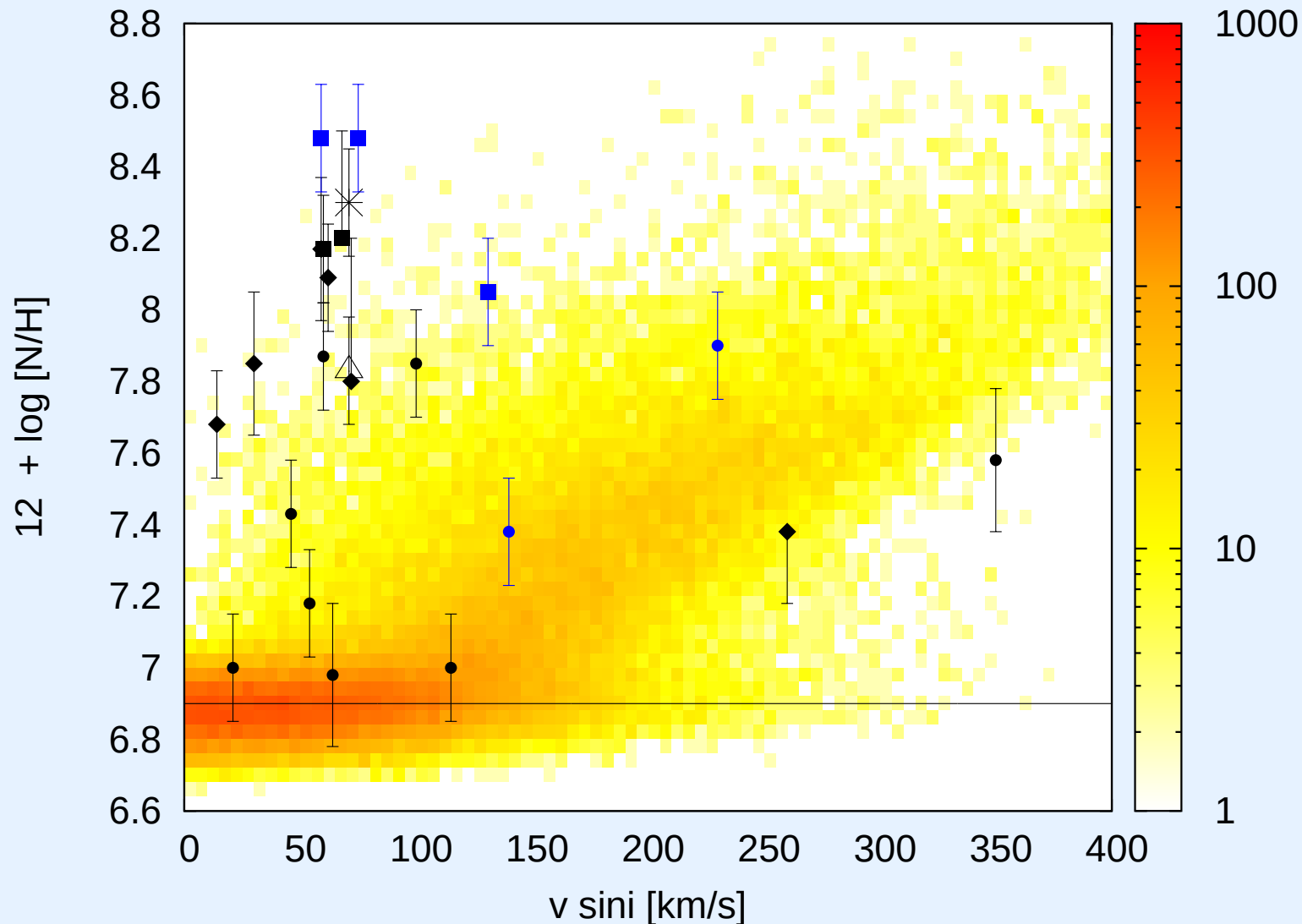
The Hunter diagram: early BV stars



The rotation rates: early BV stars

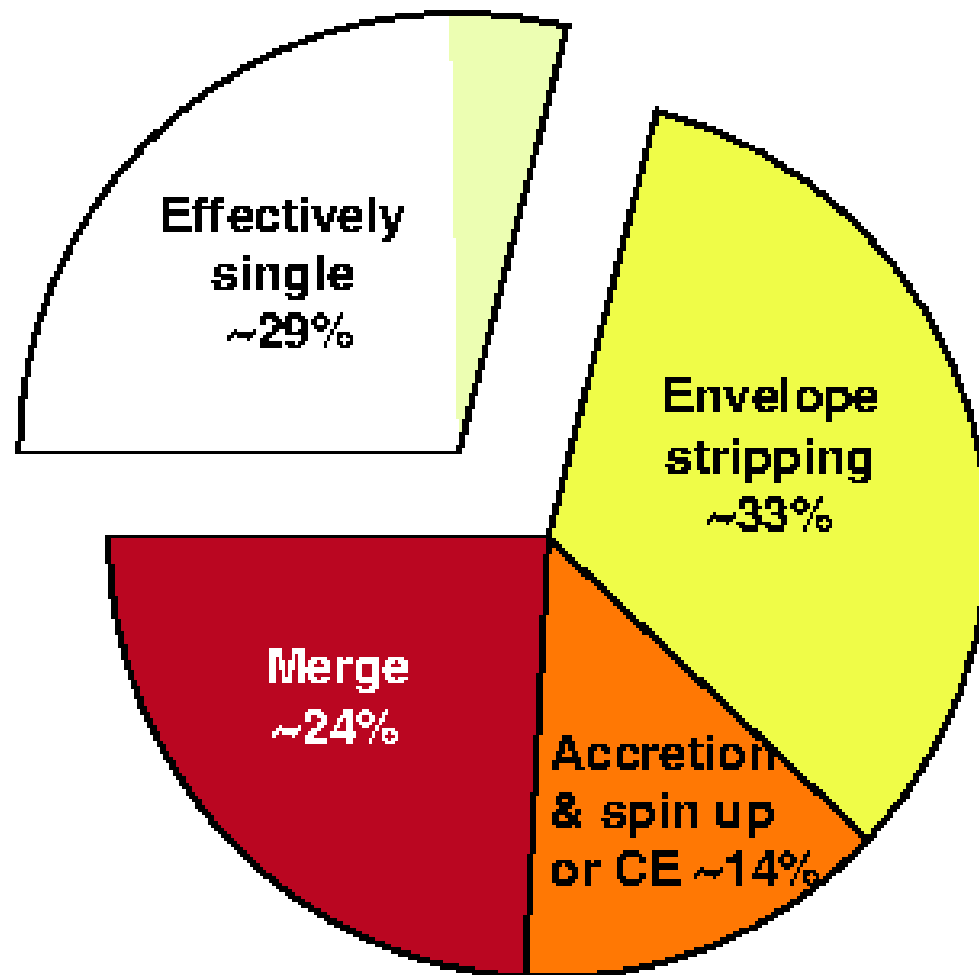


The Hunter diagram: OV stars

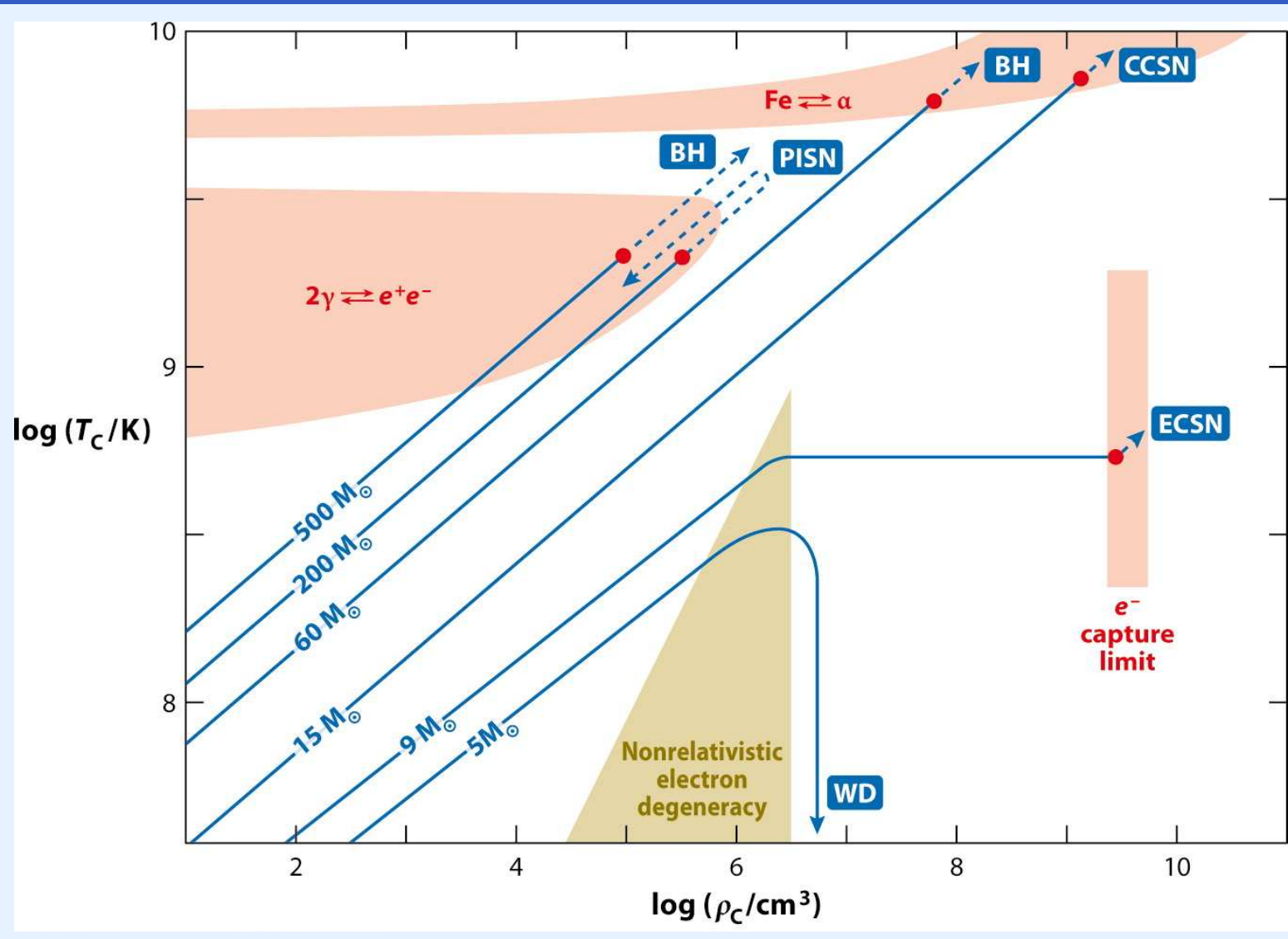


Rivero Gonzalez, Puls, Najarro & Brott 2012

Close binaries: O stars



Why mass matters



The Eddington limit

• $1 = \frac{L}{L_{\text{Edd}}} = \frac{1}{4\pi cG} \frac{\kappa_e L}{M} :$ no upper mass limit

The Eddington limit

• $1 = \frac{L}{L_{\text{Edd}}} = \frac{1}{4\pi cG} \frac{\kappa_e L}{M} :$ no upper mass limit

• $1 = \frac{1}{4\pi cG} \frac{\kappa(Z, X, T, \rho) L}{M}$

The Eddington limit

• $1 = \frac{L}{L_{\text{Edd}}} = \frac{1}{4\pi cG} \frac{\kappa_e L}{M} : \quad \text{no upper mass limit}$

• $1 = \frac{1}{4\pi cG} \frac{\kappa(Z, X, T, \rho) L}{M}$

• $1 - \left(\frac{v_{\text{rot}}^2}{v_{\text{Kepler}}^2} \right) = \frac{1}{4\pi cG} \frac{\kappa(Z, X, T, \rho) L}{M} \rightarrow \text{low } L; \text{ 2D}$

The Eddington limit

- $1 = \frac{L}{L_{\text{Edd}}} = \frac{1}{4\pi cG} \frac{\kappa_e L}{M}$: no upper mass limit

- $1 = \frac{1}{4\pi cG} \frac{\kappa(Z, X, T, \rho) L}{M}$

- $1 - \left(\frac{v_{\text{rot}}^2}{v_{\text{Kepler}}^2} \right) = \frac{1}{4\pi cG} \frac{\kappa(Z, X, T, \rho) L}{M} \rightarrow \text{low } L; \text{ 2D}$

- $1 - \left(\frac{v_{\text{rot}}^2}{v_{\text{Kepler}}^2} \right) = \frac{1}{4\pi cG} \frac{\kappa(r)(L(r) - L_{\text{conv}}(r))}{M(r)}$

The Eddington limit

- $1 = \frac{L}{L_{\text{Edd}}} = \frac{1}{4\pi cG} \frac{\kappa_e L}{M}$: no upper mass limit

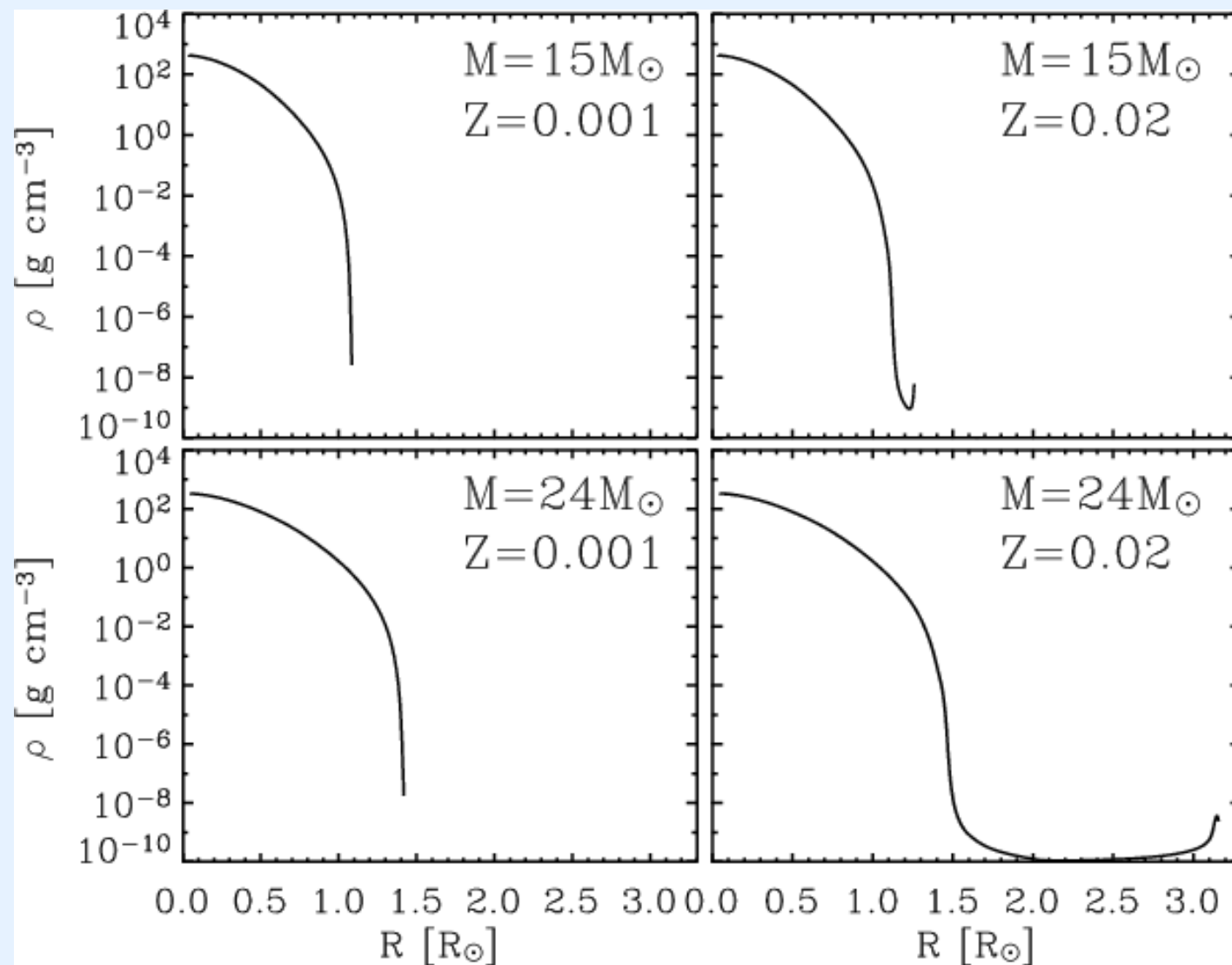
- $1 = \frac{1}{4\pi cG} \frac{\kappa(Z, X, T, \rho) L}{M}$

- $1 - \left(\frac{v_{\text{rot}}^2}{v_{\text{Kepler}}^2} \right) = \frac{1}{4\pi cG} \frac{\kappa(Z, X, T, \rho) L}{M} \rightarrow \text{low } L; \text{ 2D}$

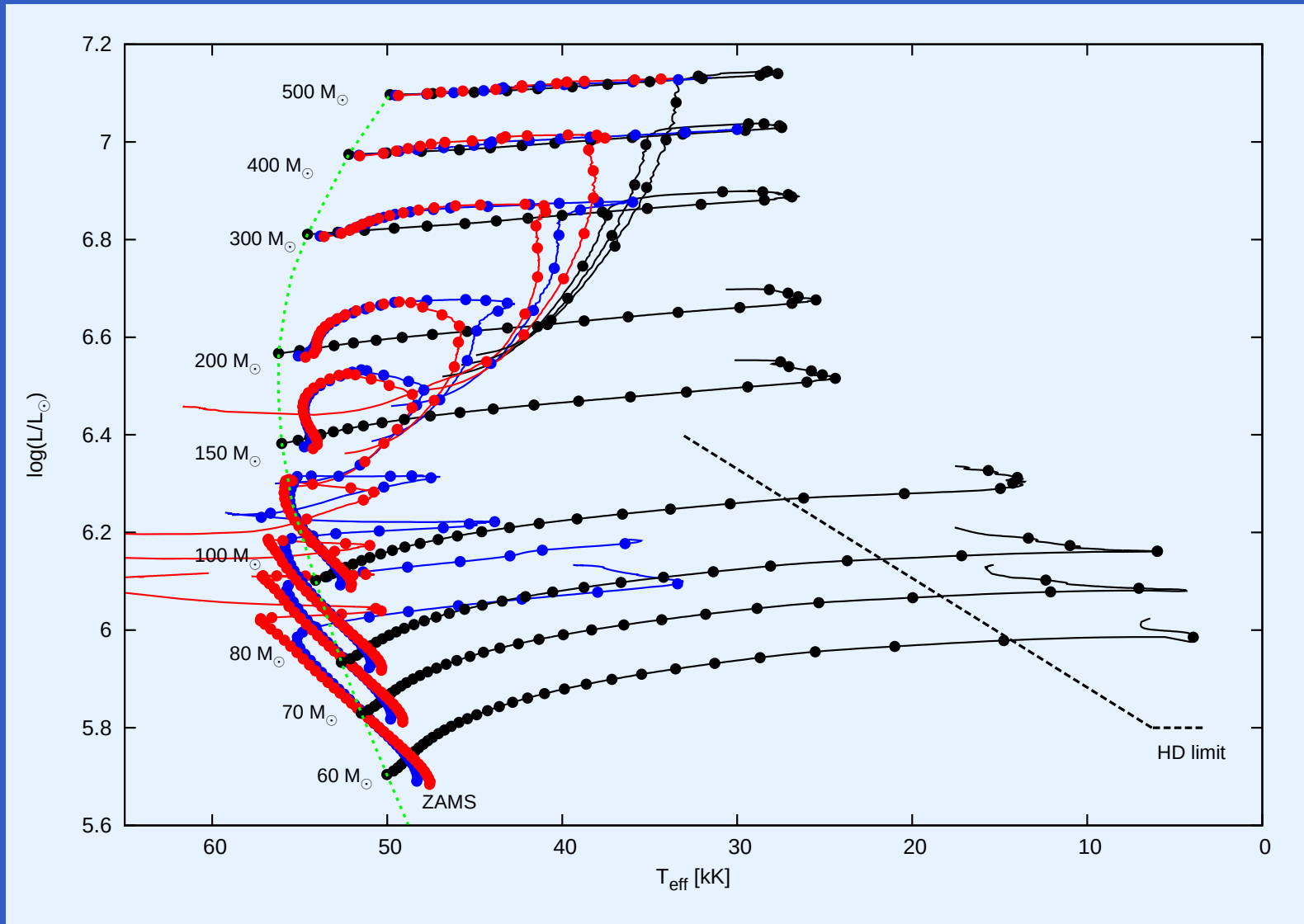
- $1 - \left(\frac{v_{\text{rot}}^2}{v_{\text{Kepler}}^2} \right) = \frac{1}{4\pi cG} \frac{\kappa(r)(L(r) - L_{\text{conv}}(r))}{M(r)}$

- inflation!

Inflation: He star models



Inflation $\neq$ HD-limit



Inflation of main sequence stars

- inflation as $f(Z)$:

H-ZAMS: $Z_{\odot} \rightarrow M < 120M_{\odot}$ Ishii et al. 1999

$Z_{SMC} \rightarrow M < 200M_{\odot}$

$Z = 0 \rightarrow M < 1000M_{\odot}$ Yoon et al. 2012

He-ZAMS: $Z_{\odot} \rightarrow M < 15M_{\odot}$ Petrovic et al. 2006

$Z_{SMC} \rightarrow M < 25M_{\odot}$

Inflation of main sequence stars

- inflation as $f(Z)$:

$$\text{H-ZAMS: } Z_{\odot} \rightarrow M < 120 M_{\odot} \quad \text{Ishii et al. 1999}$$

$$Z_{SMC} \rightarrow M < 200 M_{\odot}$$

$$Z = 0 \rightarrow M < 1000 M_{\odot} \quad \text{Yoon et al. 2012}$$

$$\text{He-ZAMS: } Z_{\odot} \rightarrow M < 15 M_{\odot} \quad \text{Petrovic et al. 2006}$$

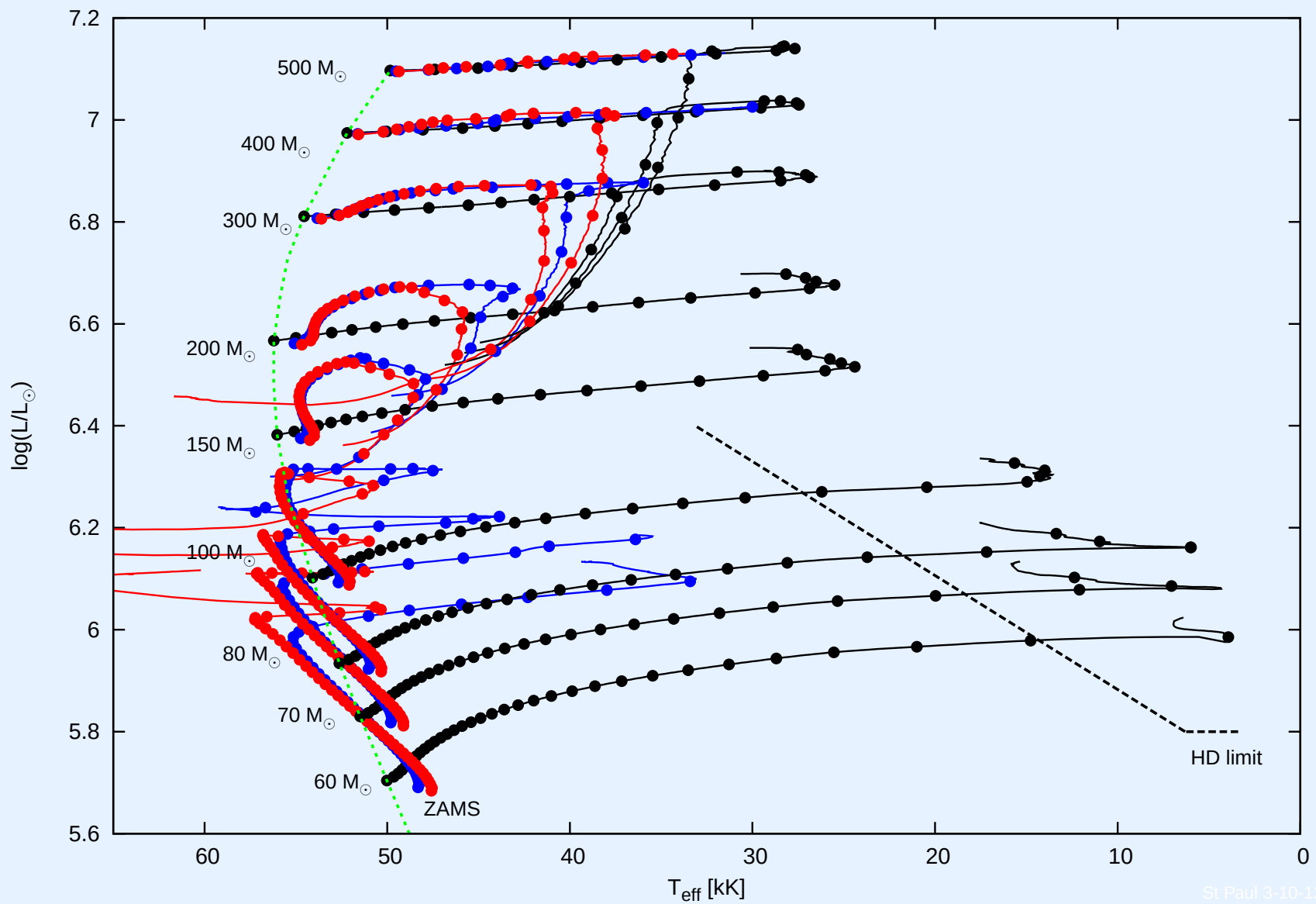
$$Z_{SMC} \rightarrow M < 25 M_{\odot}$$

- $\frac{L}{L_{\text{Edd}}} \sim \frac{\kappa L}{M}$:

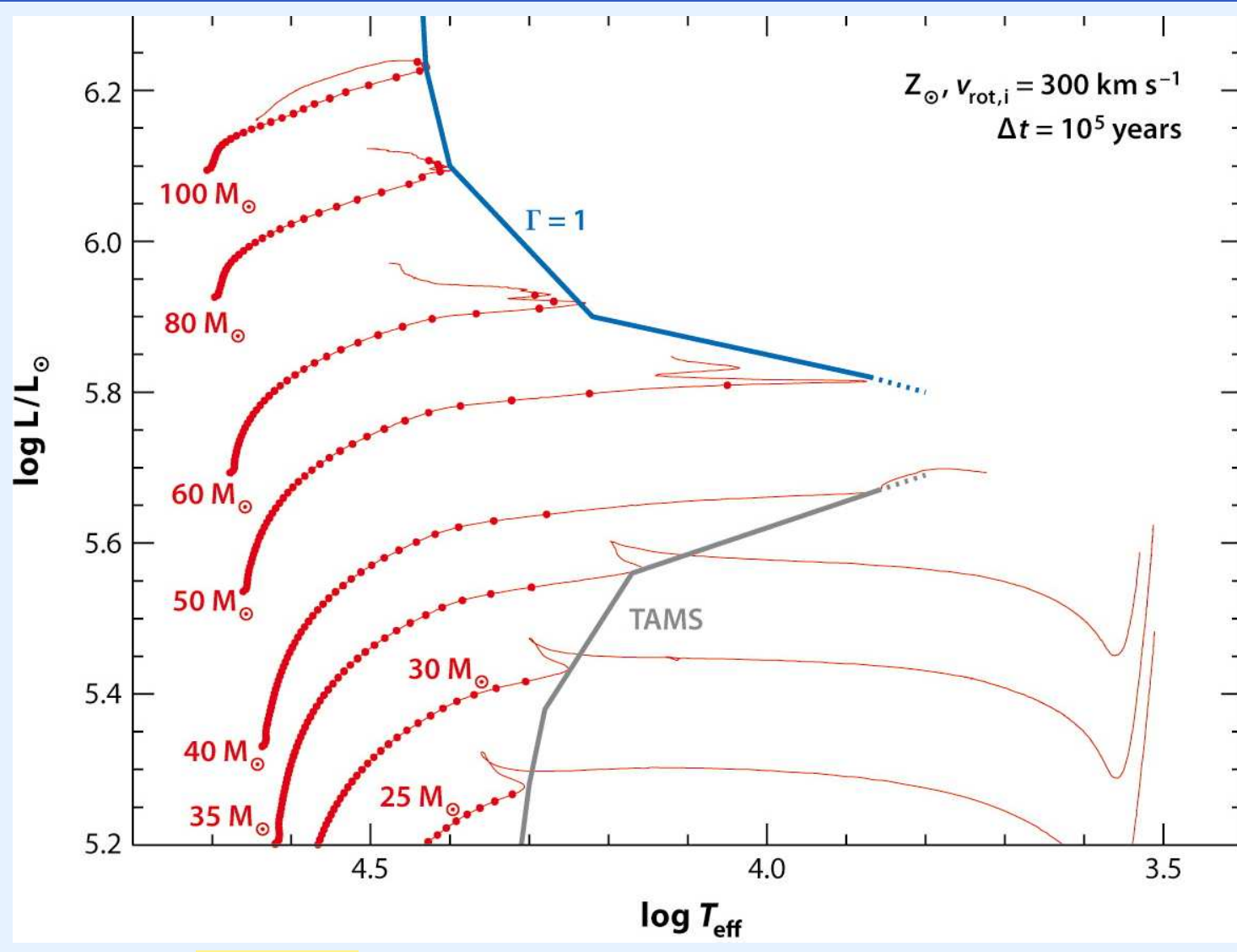
$$\text{H} \rightarrow \text{He} \Rightarrow \kappa \downarrow 2, L(M) \uparrow 10$$

$\Rightarrow$ He-ZAMS crosses H-ZAMS!

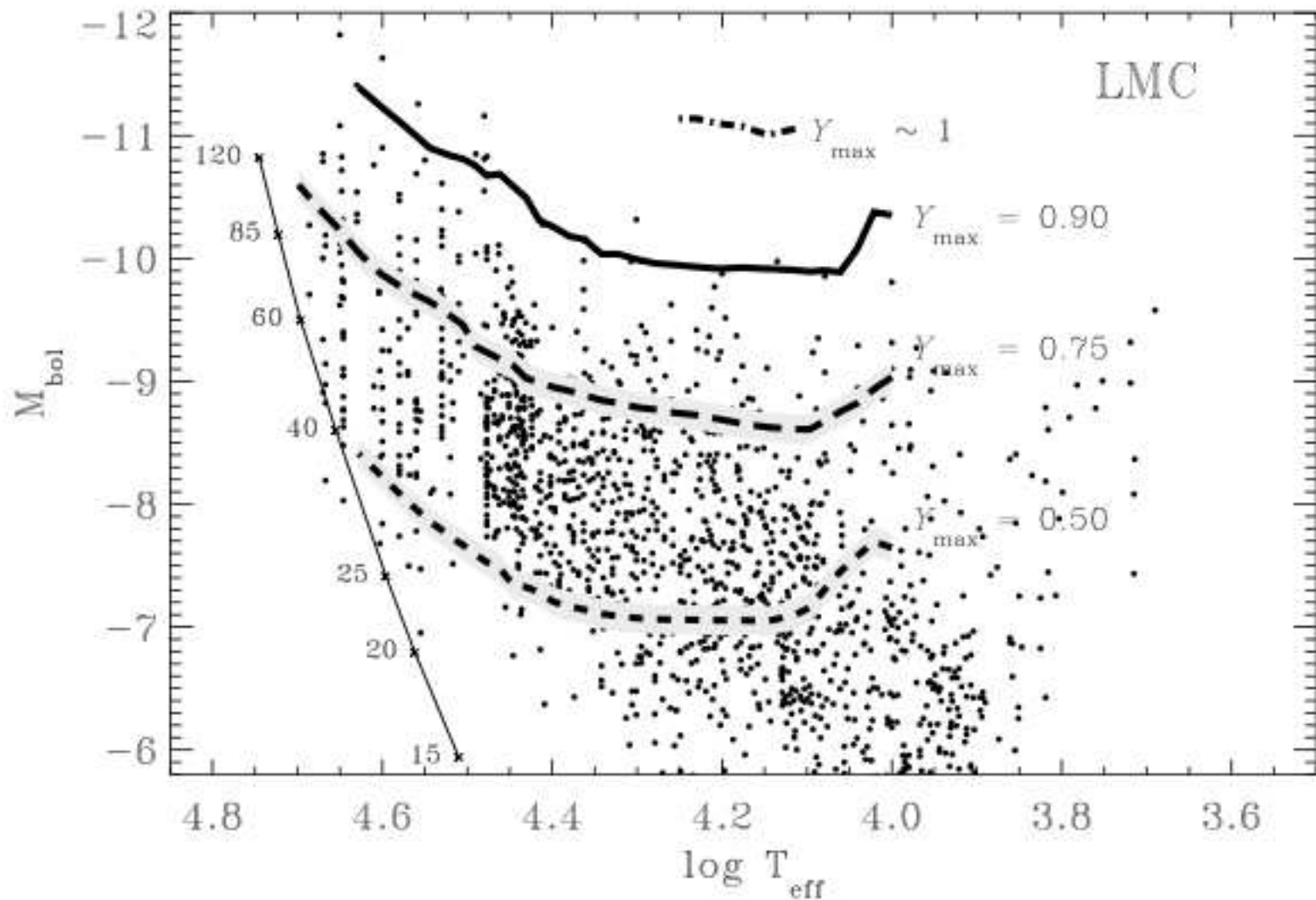
ZAMS crossing



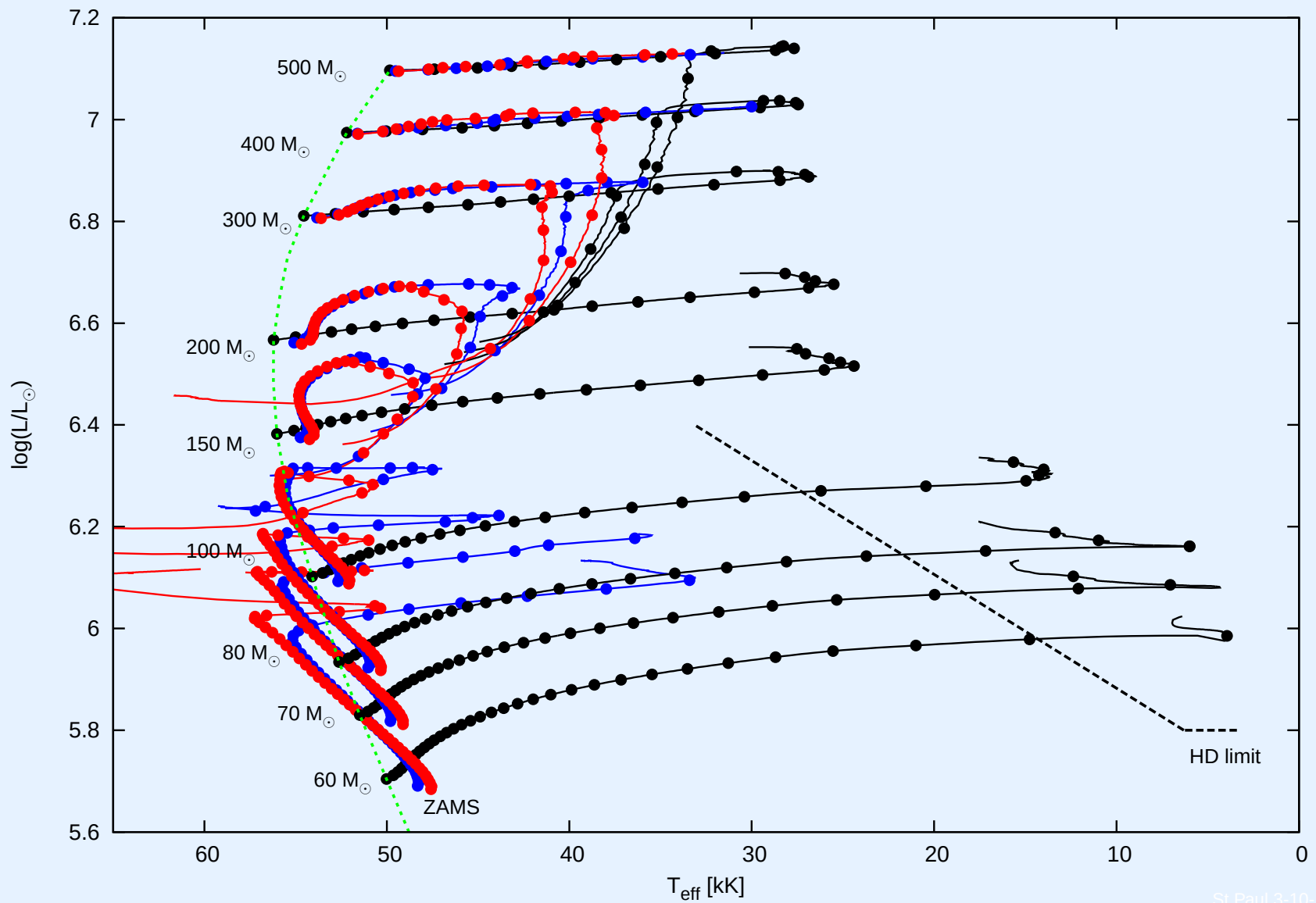
Eddington limit



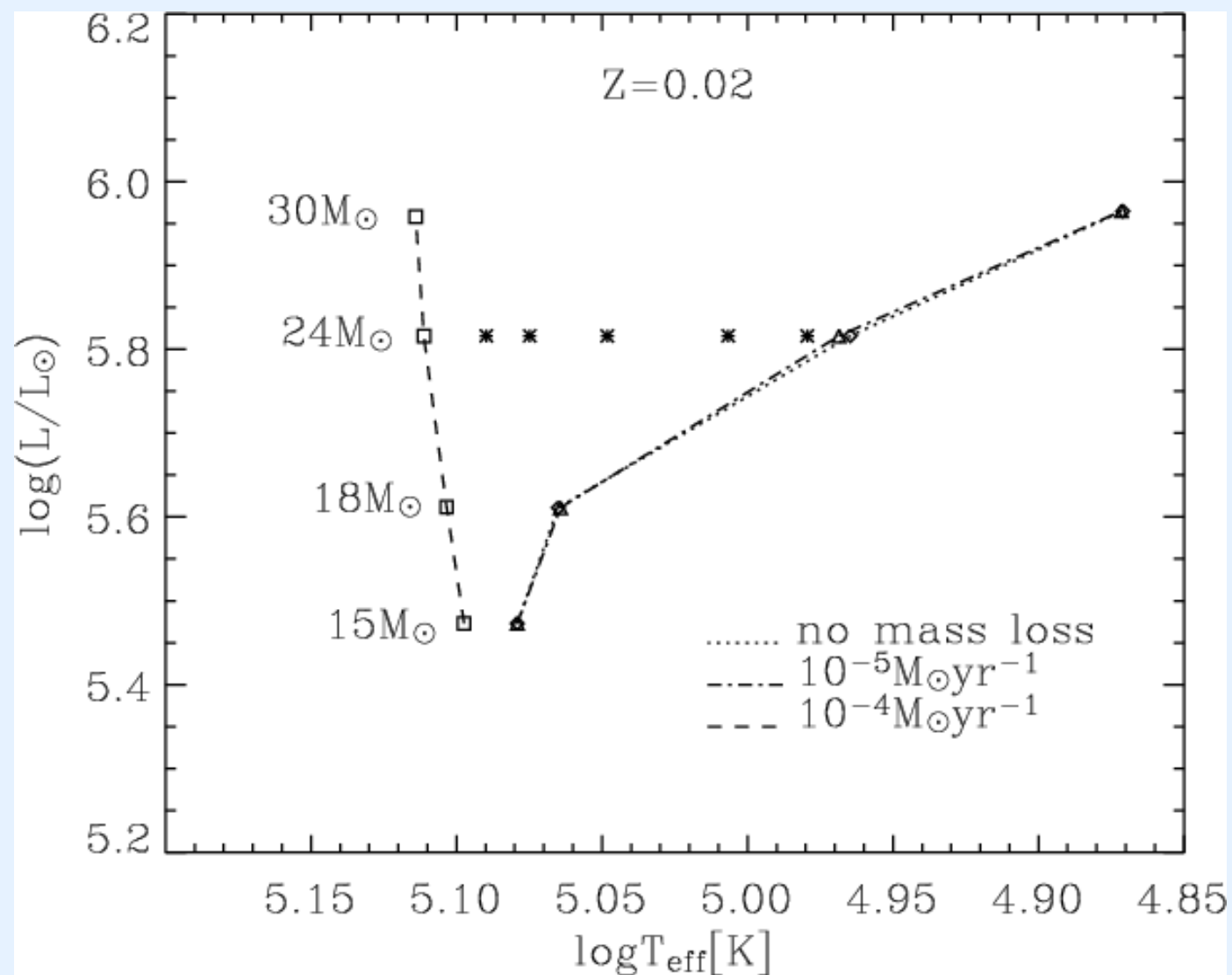
Eddington limit



Eddington-limit



Inflation: $f(\dot{M})$



Inflation: S Dor Variations?

Ingredients:

- $R(\text{inflated layer}) = f(\dot{M})$

Petrovic et al. 2006

- $\dot{M} = f(R)$

e.g. Vink: bistability

Inflation: S Dor Variations?

Ingredients:

• $R(\text{inflated layer}) = f(\dot{M})$

Petrovic et al. 2006

• $\dot{M} = f(R)$

e.g. Vink: bistability

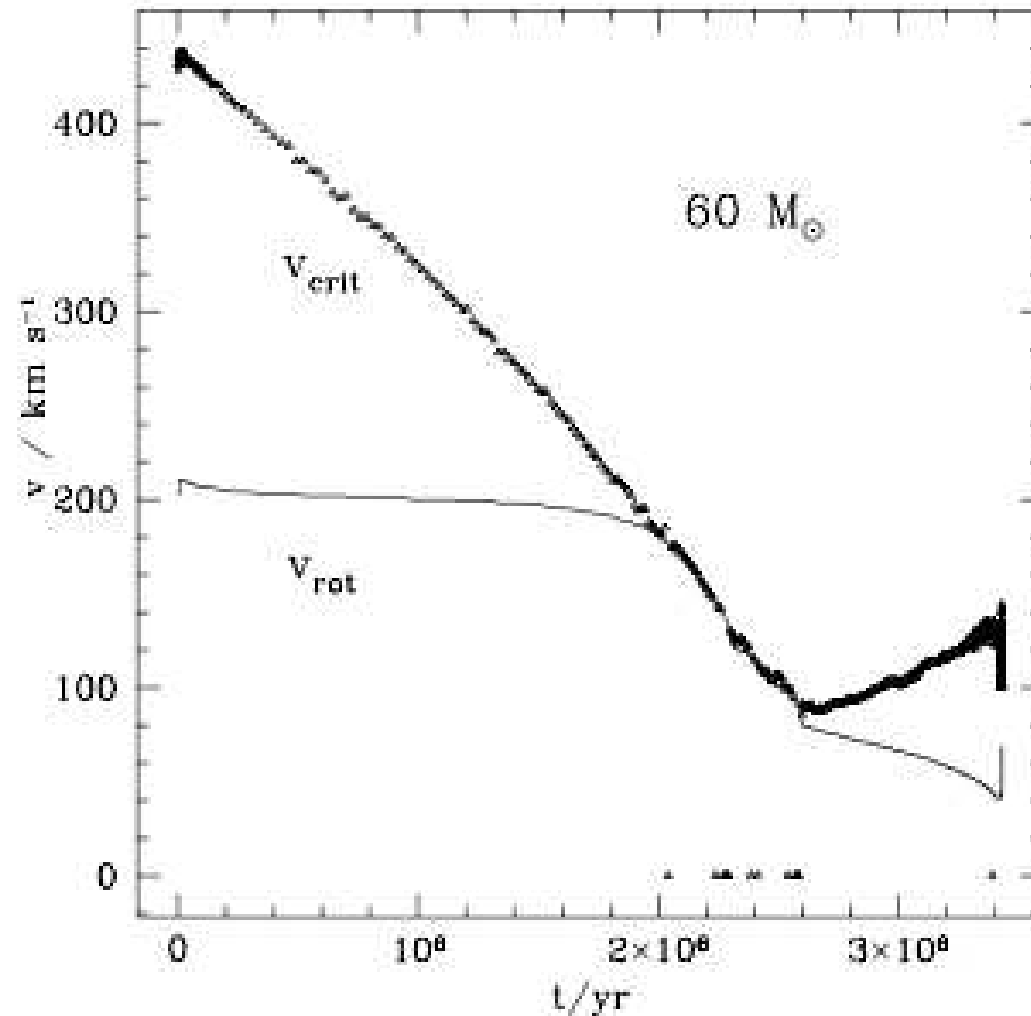
1) star compact: $\dot{M}$ small

→ 2) inflation

→ 3) $\dot{M}$ large

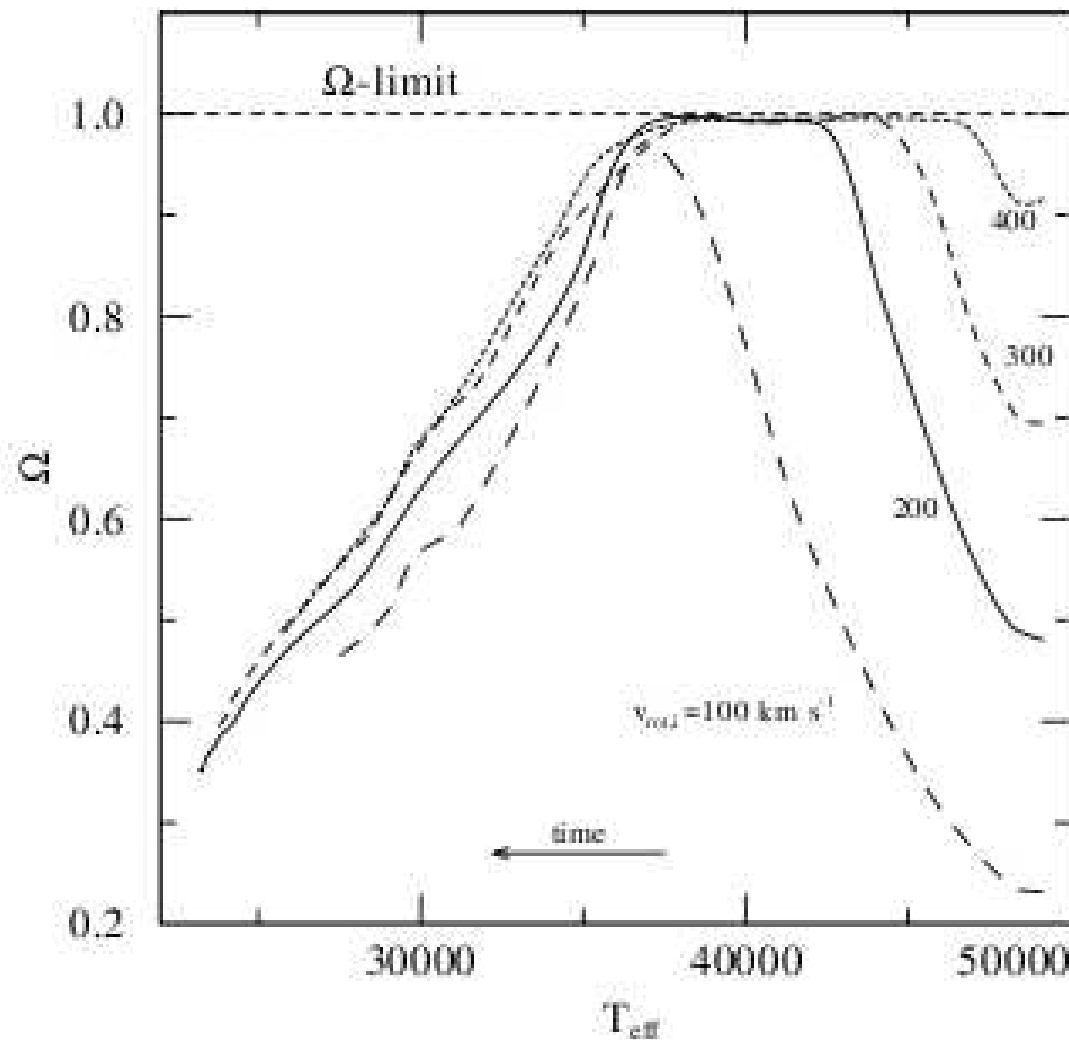
→ 1) star compact

The role of stellar evolution



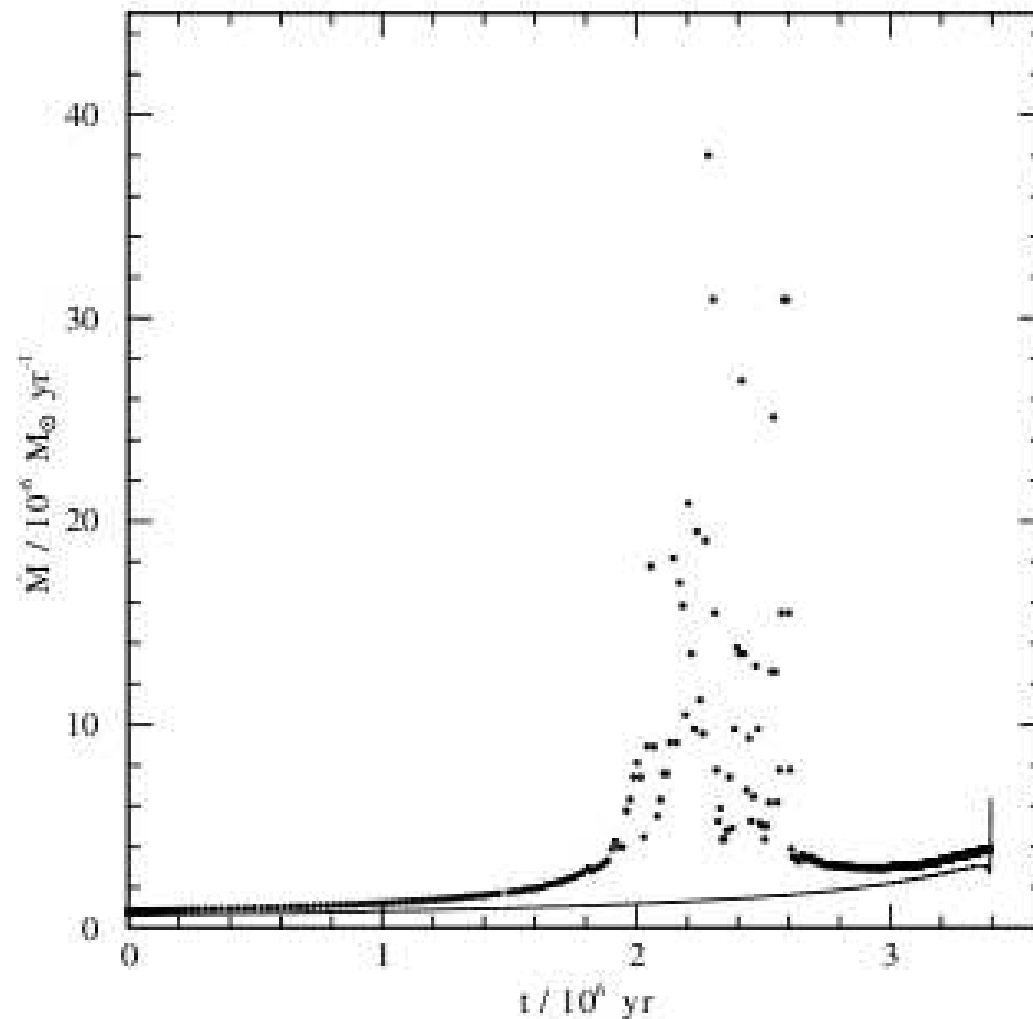
Langer 1998

MS evolution & the Ω -limit



Langer 1998

MS evolution: $\dot{M}$



Langer 1998

Mass loss rate at the Edd.-limit

nuclear timescale evolution: $\dot{M} \simeq \frac{M}{\tau_{\text{nuc}}}$

$$M = 100 M_{\odot}: \quad \tau_{\text{nuc}} = 3 \cdot 10^6 \text{ yr} \rightarrow \dot{M} \simeq 3 \cdot 10^{-5} M_{\odot} \text{ yr}^{-1}$$

Mass loss rate at the Edd.-limit

nuclear timescale evolution: $\dot{M} \simeq \frac{M}{\tau_{\text{nuc}}}$

$$M = 100 M_{\odot}: \quad \tau_{\text{nuc}} = 3 \cdot 10^6 \text{ yr} \rightarrow \dot{M} \simeq 3 \cdot 10^{-5} M_{\odot} \text{ yr}^{-1}$$

thermal timescale evolution: $\dot{M} \simeq \frac{M}{\tau_{\text{th}}}$

$$M = 100 M_{\odot}: \quad \tau_{\text{th}} = 10^4 \text{ yr} \rightarrow \dot{M} \simeq 10^{-2} M_{\odot} \text{ yr}^{-1}$$

Mass loss rate at the Edd.-limit

nuclear timescale evolution: $\dot{M} \simeq \frac{M}{\tau_{\text{nuc}}}$

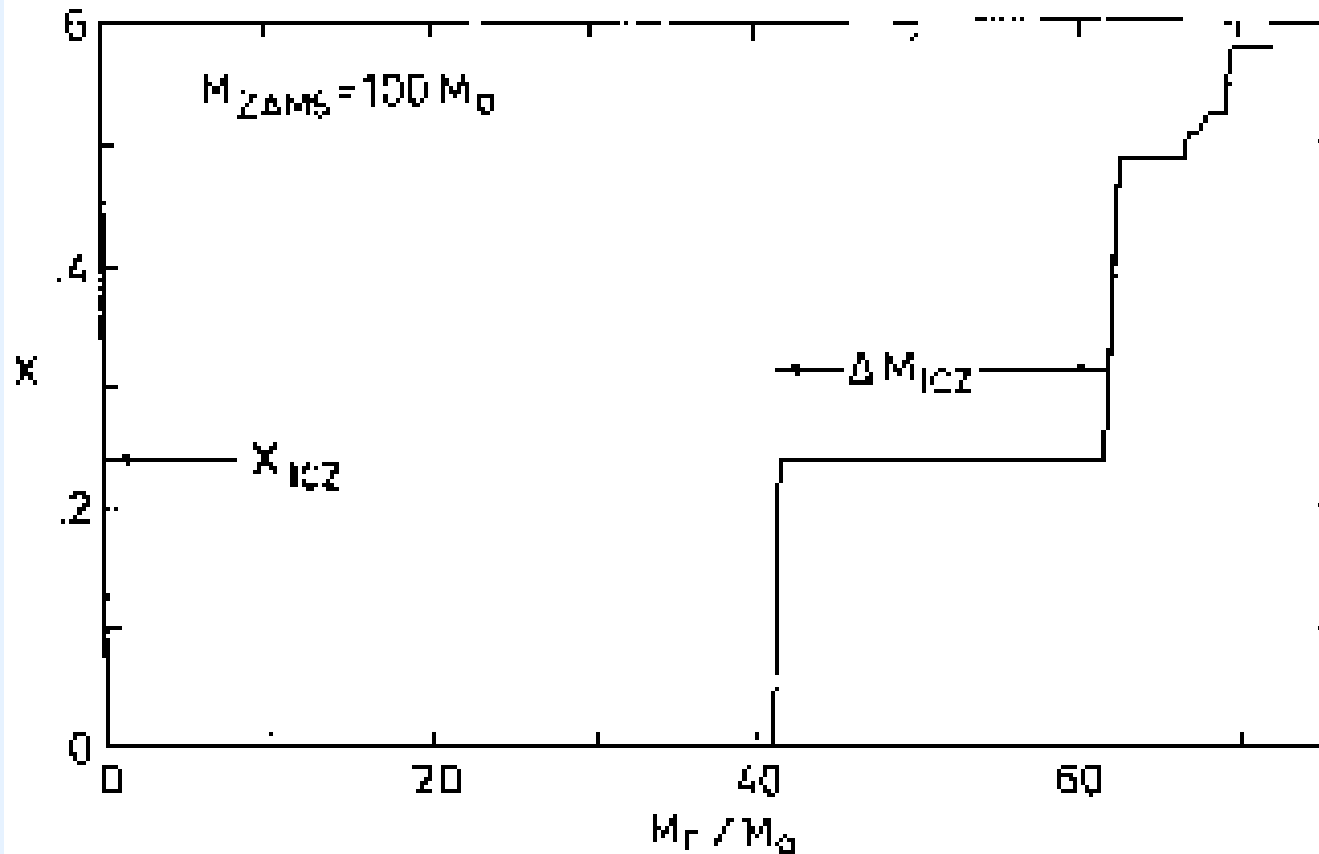
$$M = 100 M_{\odot}: \quad \tau_{\text{nuc}} = 3 \cdot 10^6 \text{ yr} \rightarrow \dot{M} \simeq 3 \cdot 10^{-5} M_{\odot} \text{ yr}^{-1}$$

thermal timescale evolution: $\dot{M} \simeq \frac{M}{\tau_{\text{th}}}$

$$M = 100 M_{\odot}: \quad \tau_{\text{th}} = 10^4 \text{ yr} \rightarrow \dot{M} \simeq 10^{-2} M_{\odot} \text{ yr}^{-1}$$

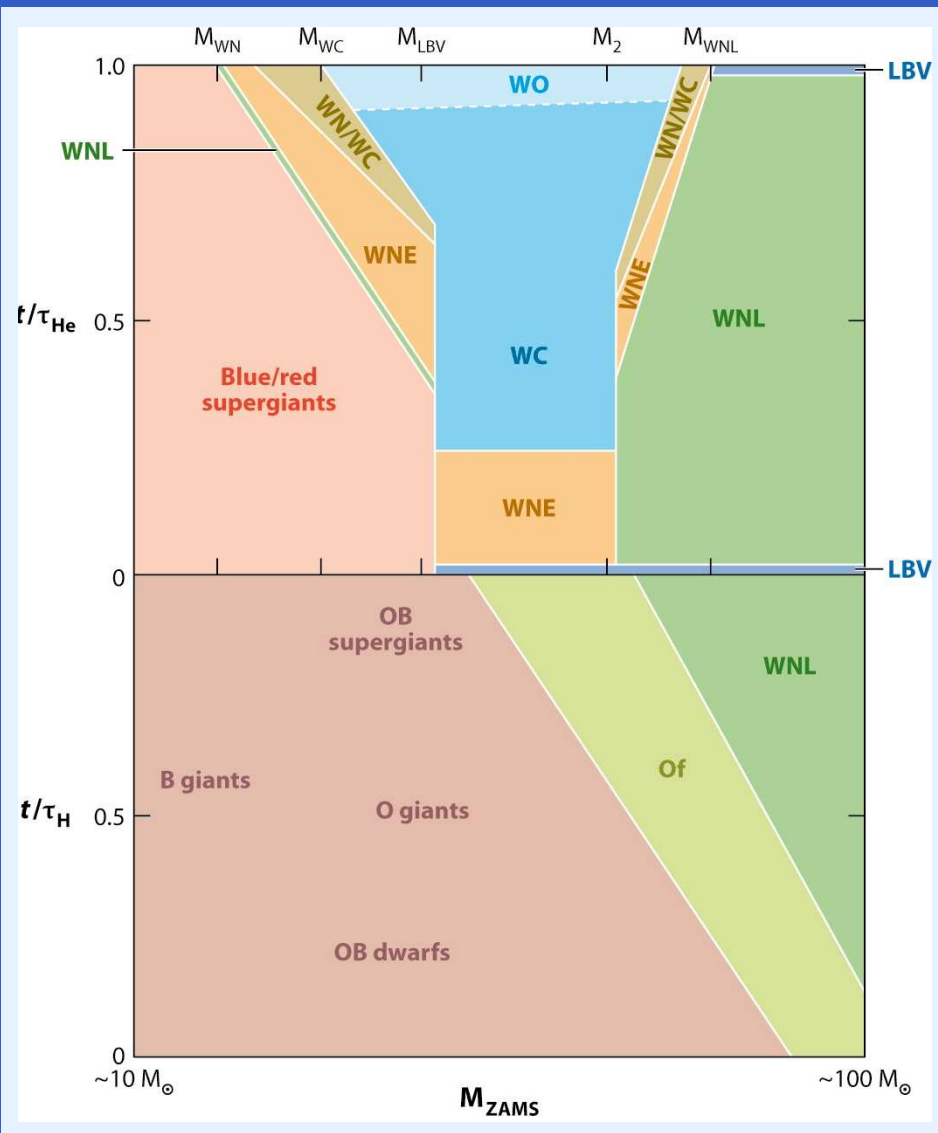
$\Rightarrow$ LBV eruptions need thermal timescale evolution

Post-LBV evolution



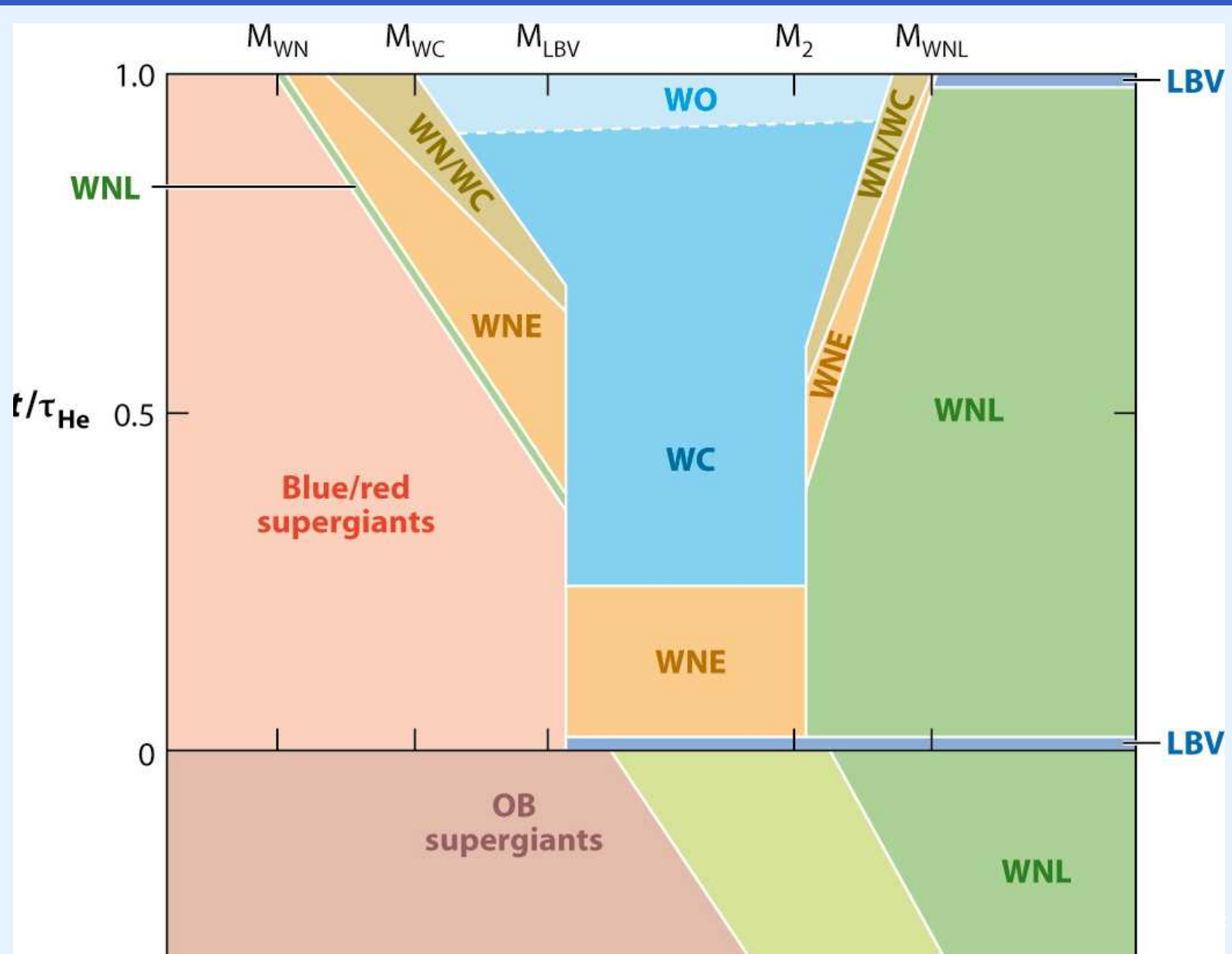
Langer 1987

Evolutionary scheme

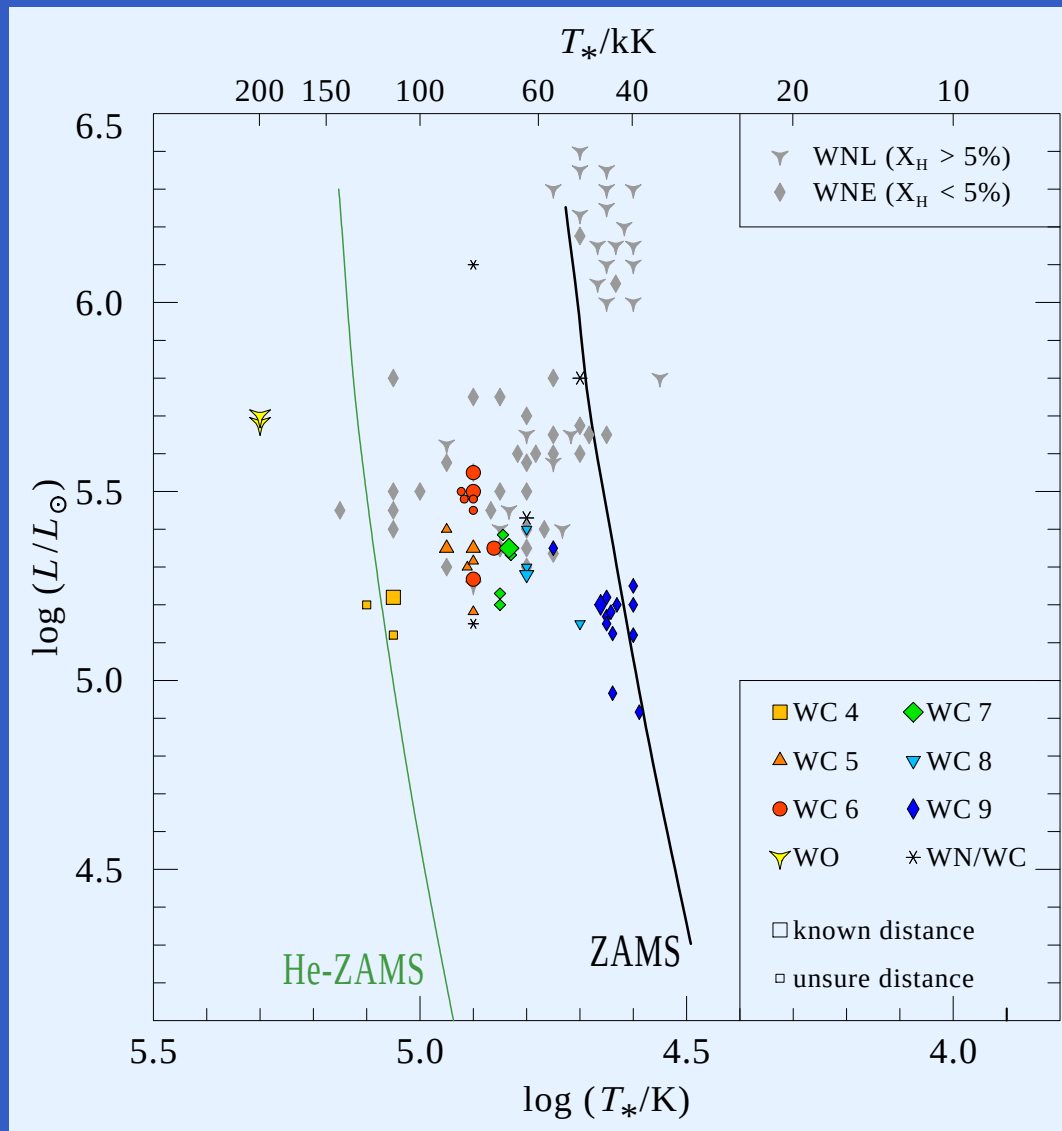


Langer 2012

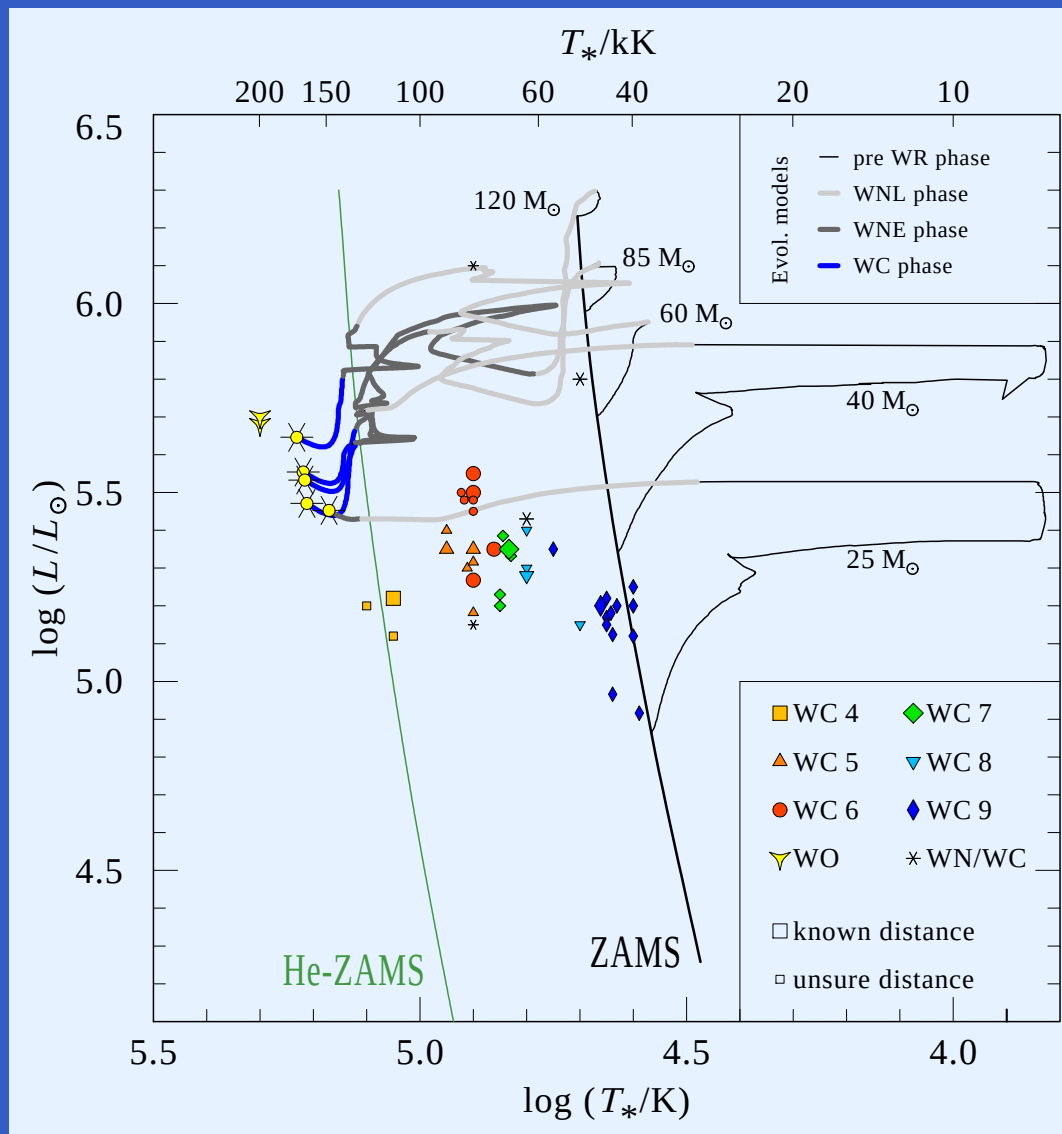
Evolutionary scheme



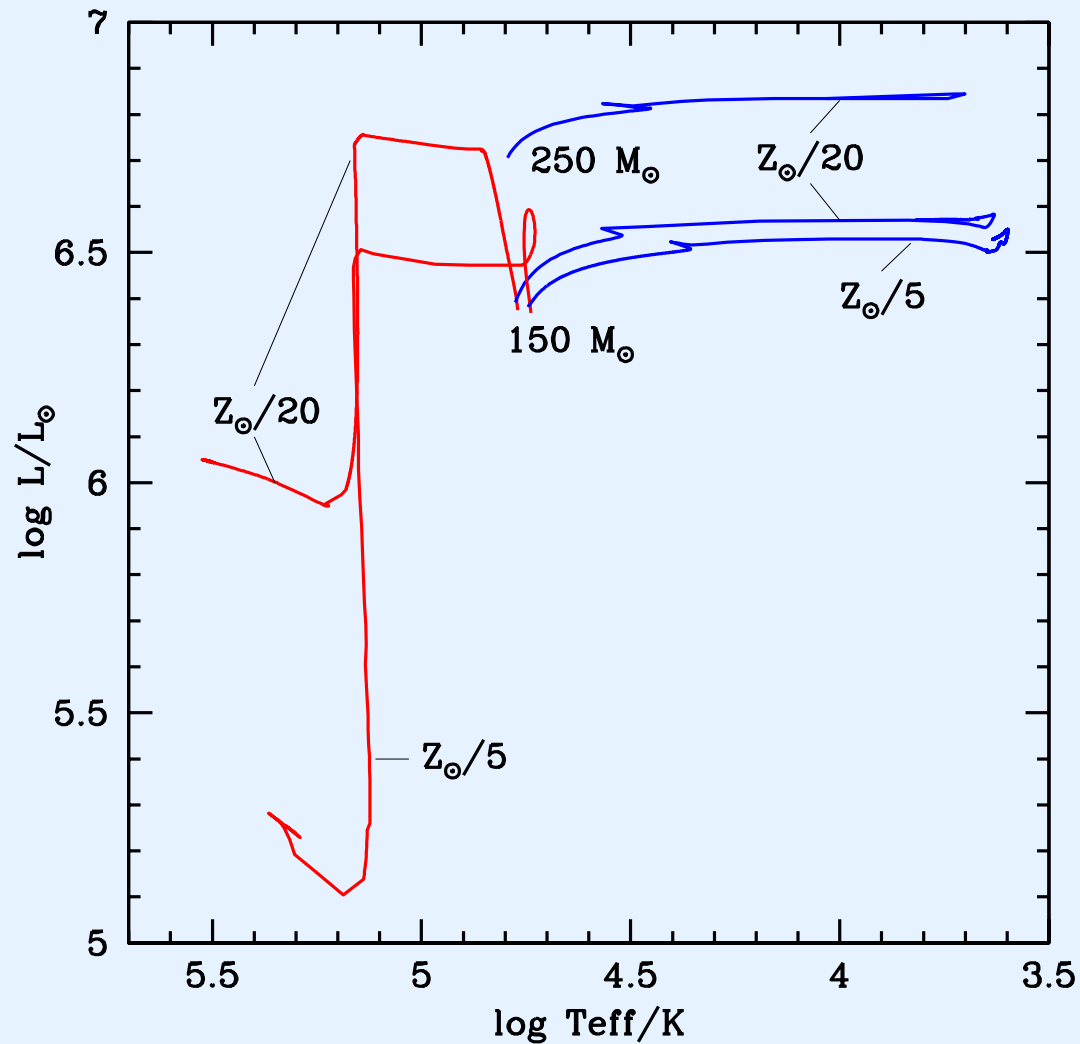
Galactic Wolf-Rayet stars



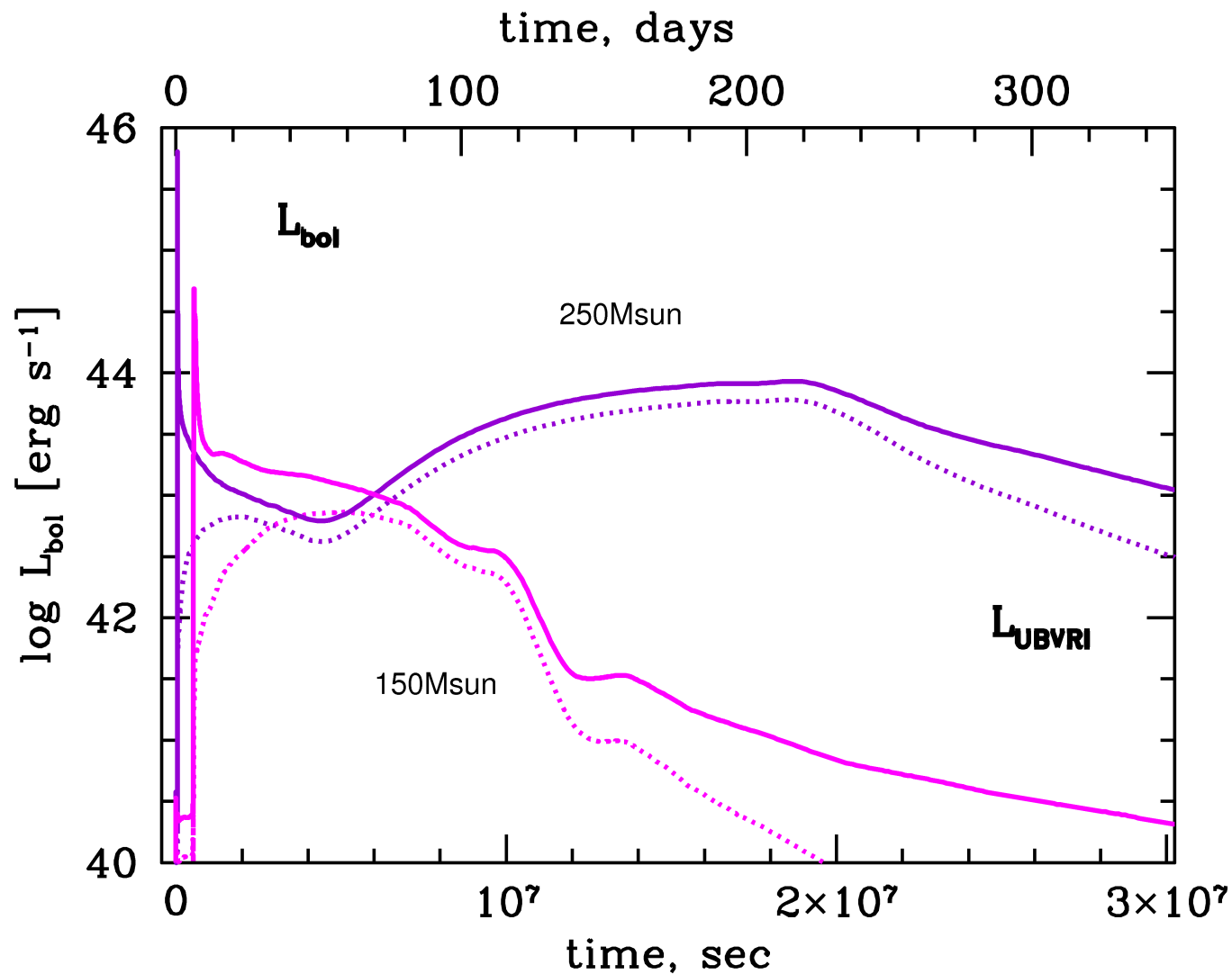
Galactic Wolf-Rayet stars



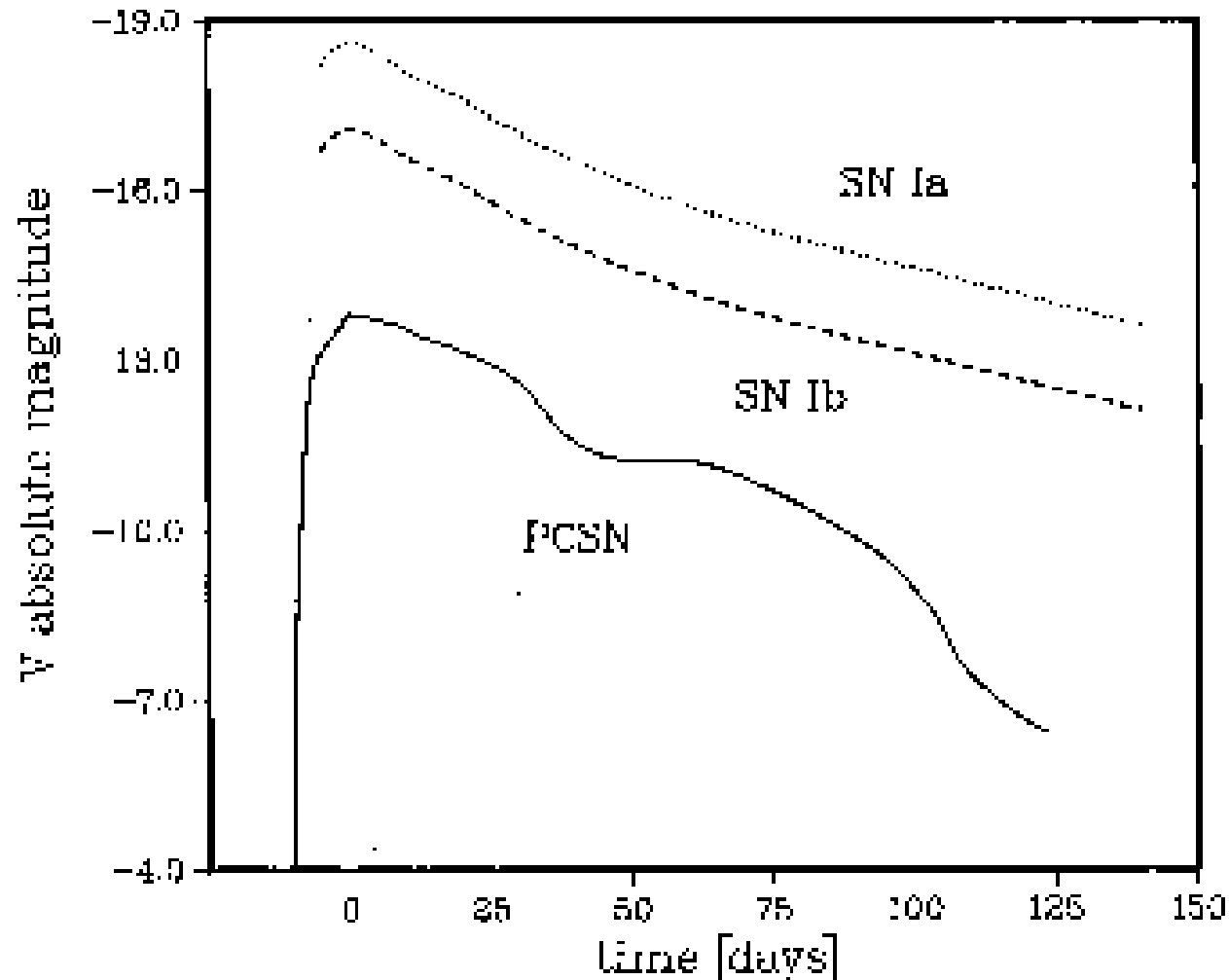
Local low-Z VMS



Local Pair-Instability SNe

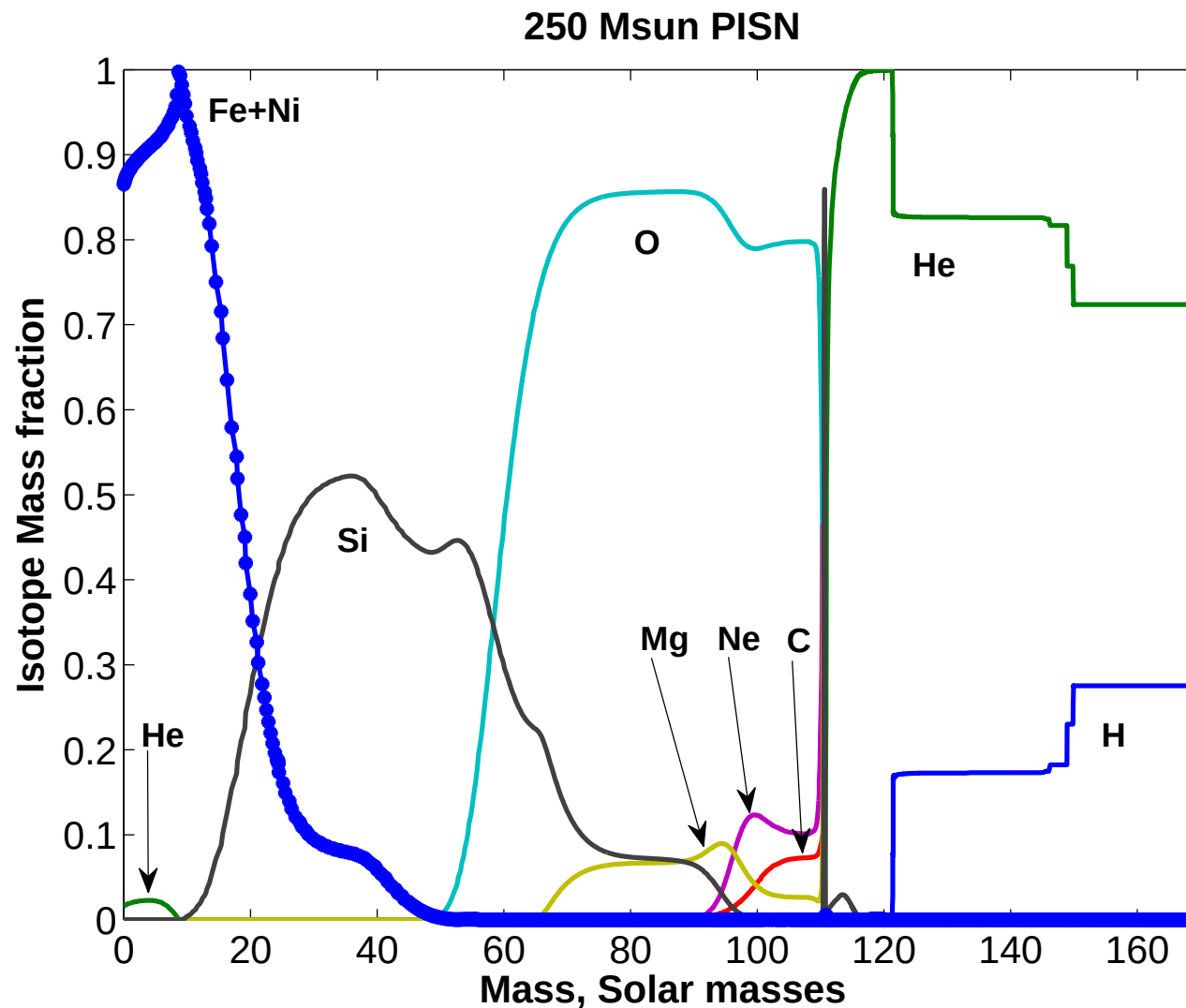


Local Pair-Instability SNe



Herzig, El Eid, Fricke & Langer 1990

Local Pair-Instability SNe: Yields



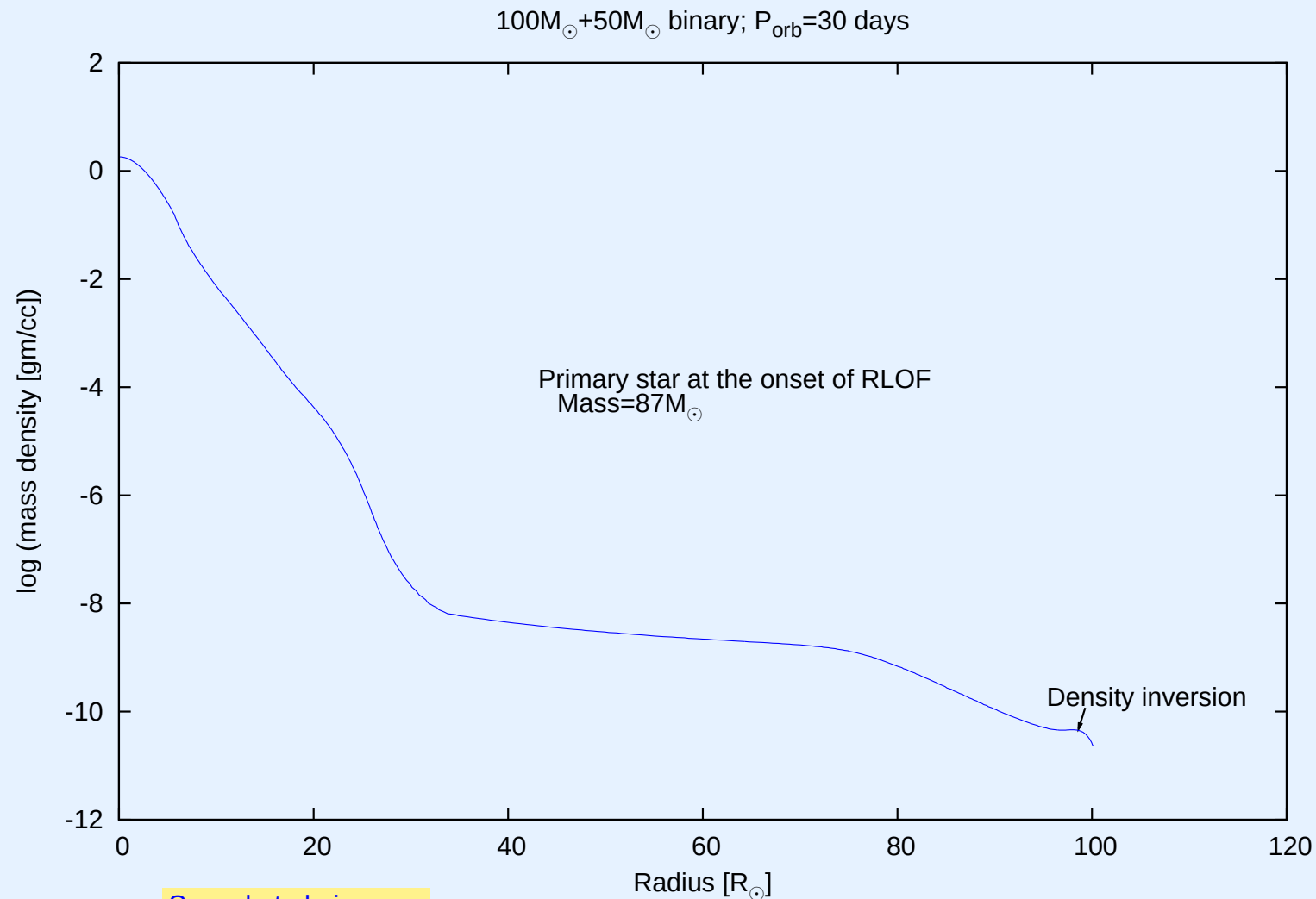
Kozyreva al., in prep.

Summary

- Complex physics at the Ω -limit
- Inflation: ZAMS crowding, S Dor variability?
- Eruptions $\leftarrow \tau_{\text{th}}$ -evolution
 $\Rightarrow$ erupting LBVs: after core-H, core-He burning
- PISNe in the local universe

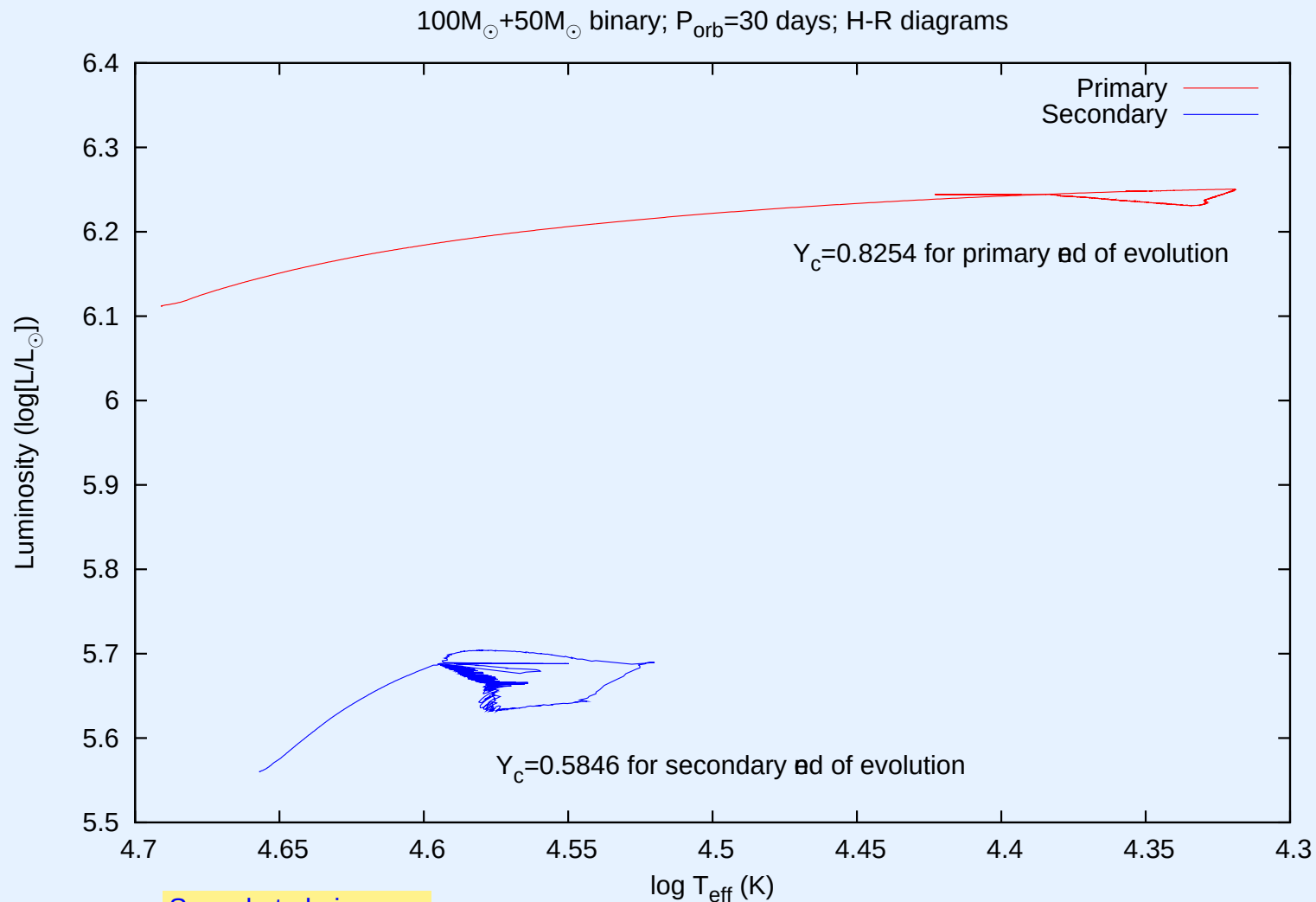
VMS binary evolution

$Y_c=0.8254$



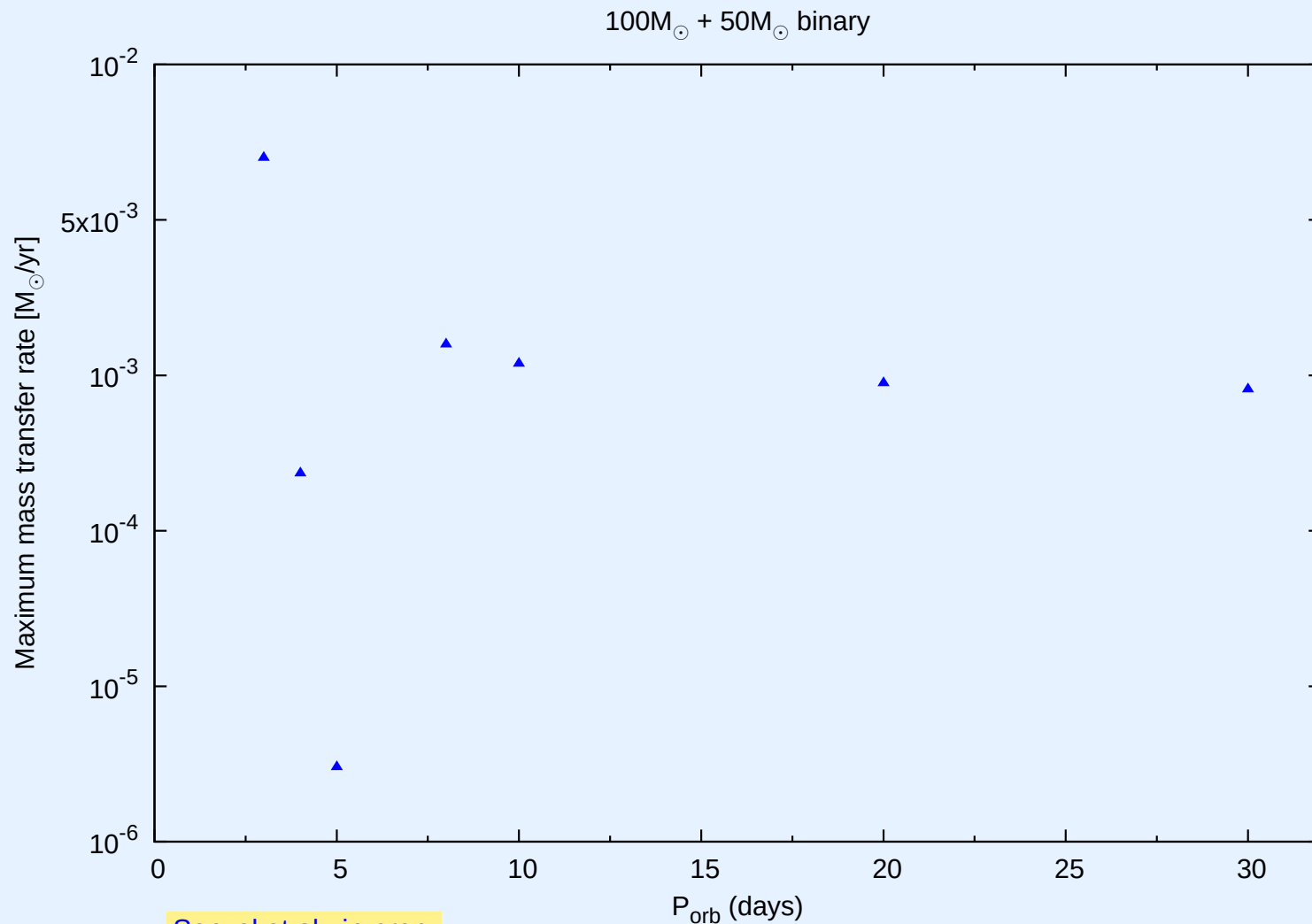
Sanyal et al., in prep.

VMS binary evolution



Sanyal et al., in prep.

VMS binary evolution



Sanyal et al., in prep.